

## Euler-Lagrange Approach to Modelling

Greg and Mario

- 1 Euler-Lagrange Equations
- 2 Modelling the Rotation Dynamics (Gros2012f, Gros2013b)
- 3 Baumgarte Stabilisation (Gros2012f)
- 4 Tether Models (Pesce2003, Zanon2012, Zanon2013a)
- 5 Kite Models (Gros2013b)

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- “*Everything Should Be Made as Simple as Possible, But Not Simpler*”, A. Einstein

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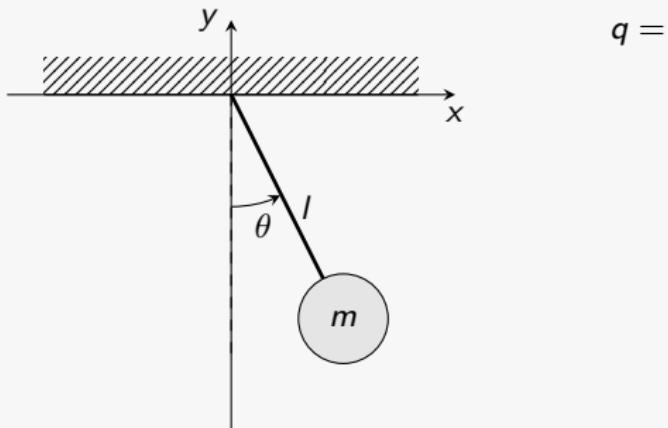
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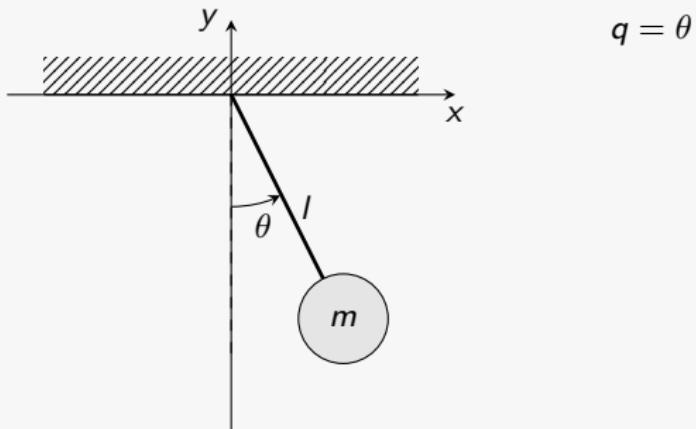
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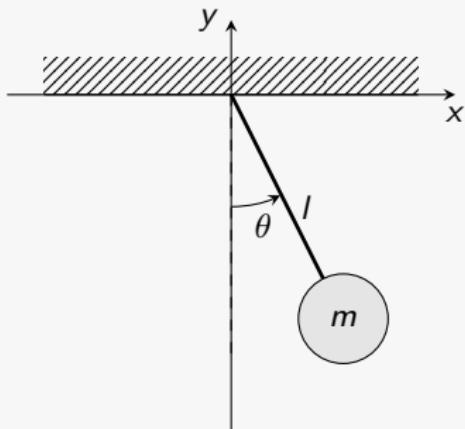
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- ⑥ Compute the generalised forces  $F^q$

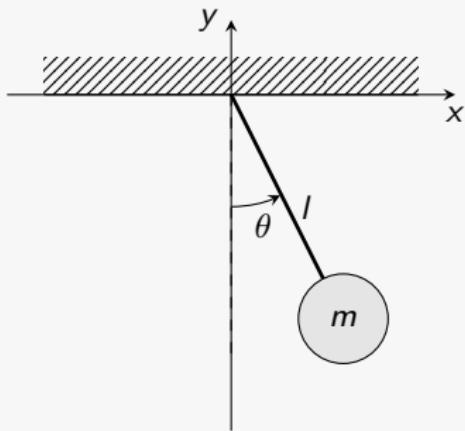
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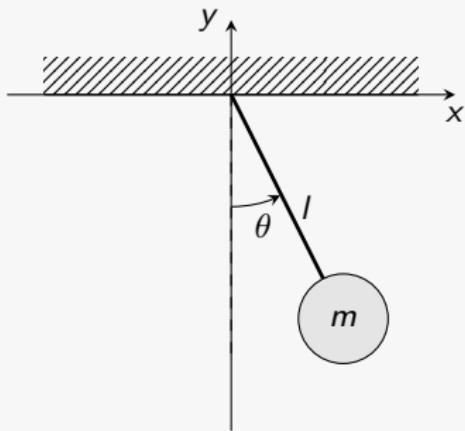
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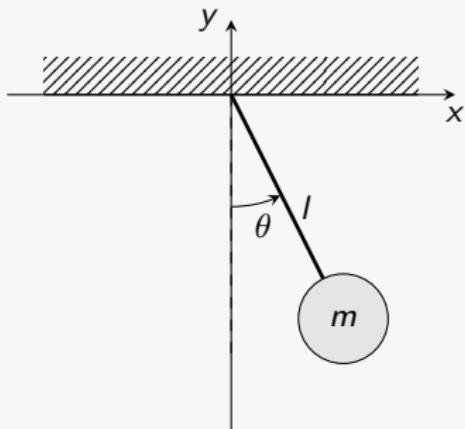
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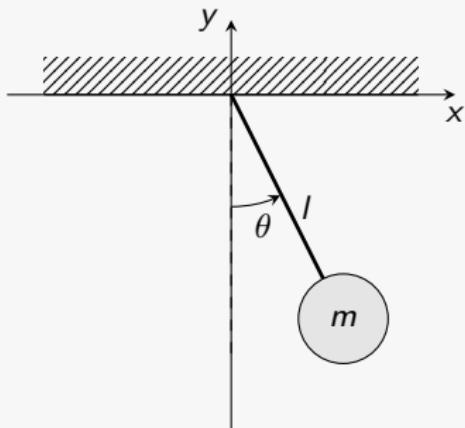
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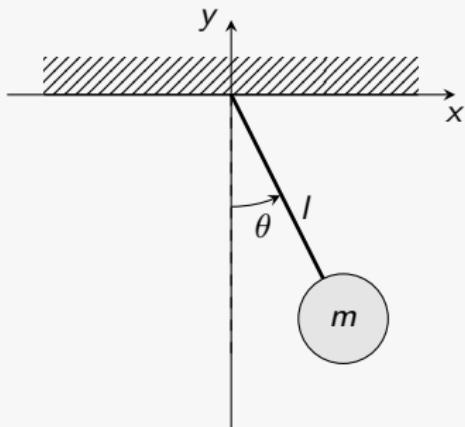
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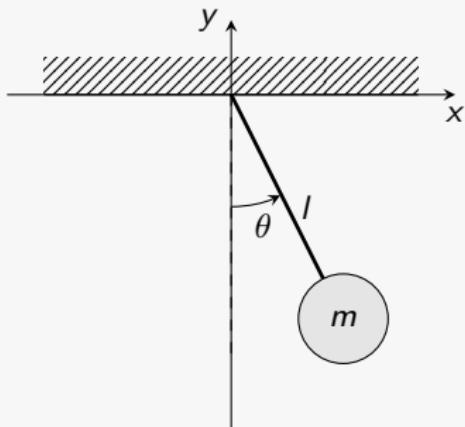
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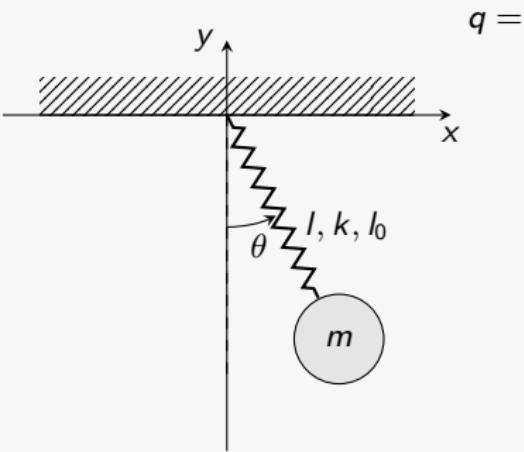
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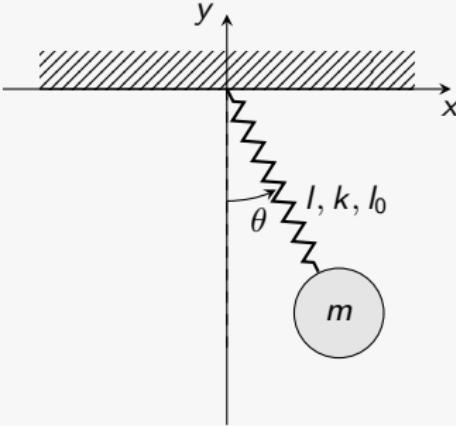
$$\ddot{\theta} = -\frac{g}{l} \sin \theta$$

## Example: Elastic Pendulum



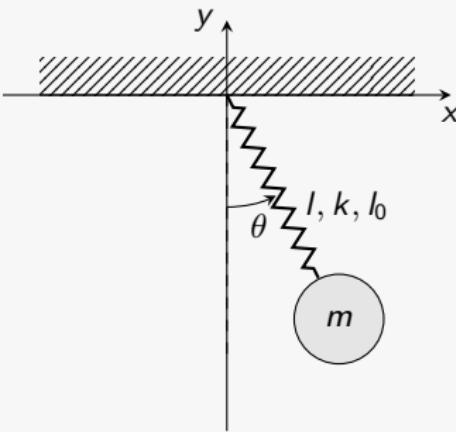
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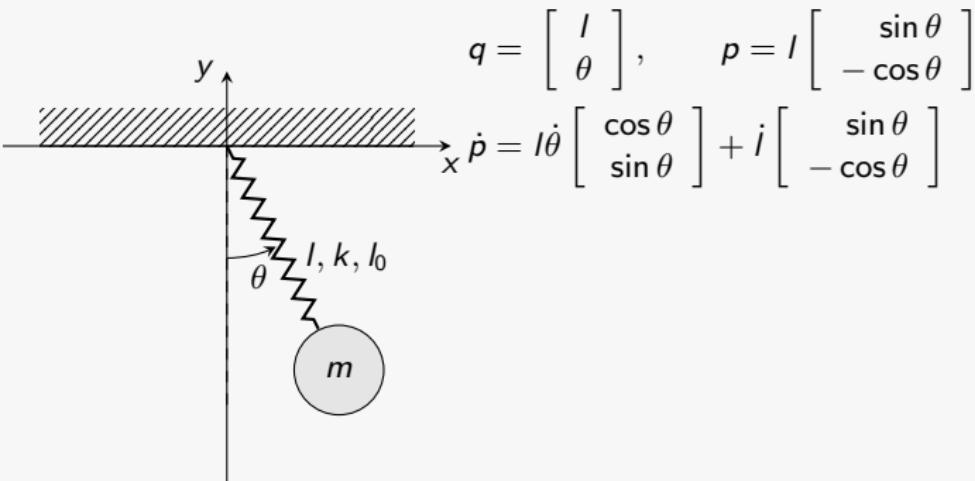


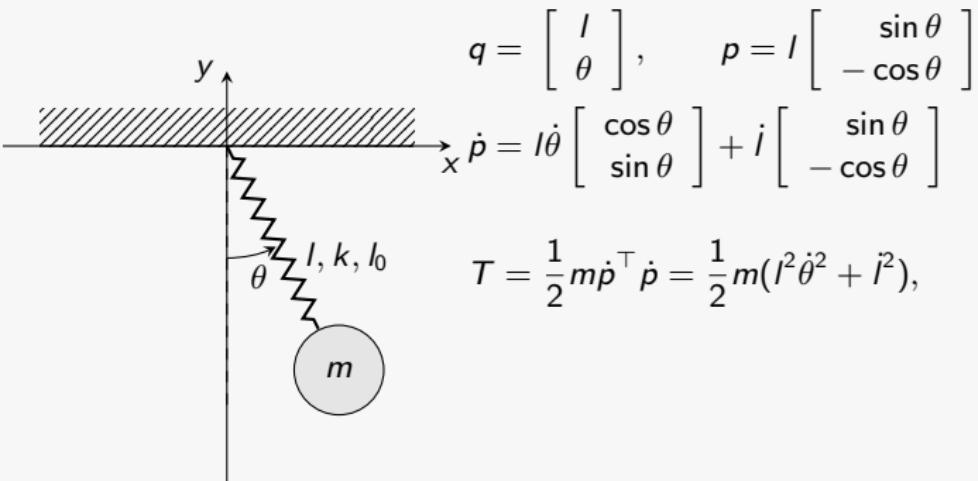
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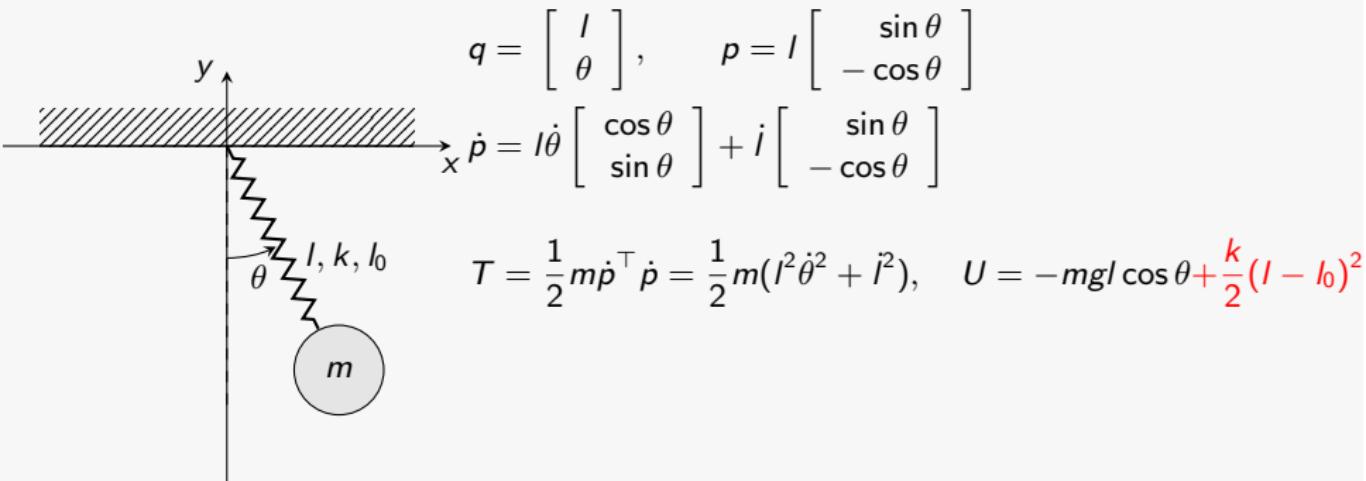


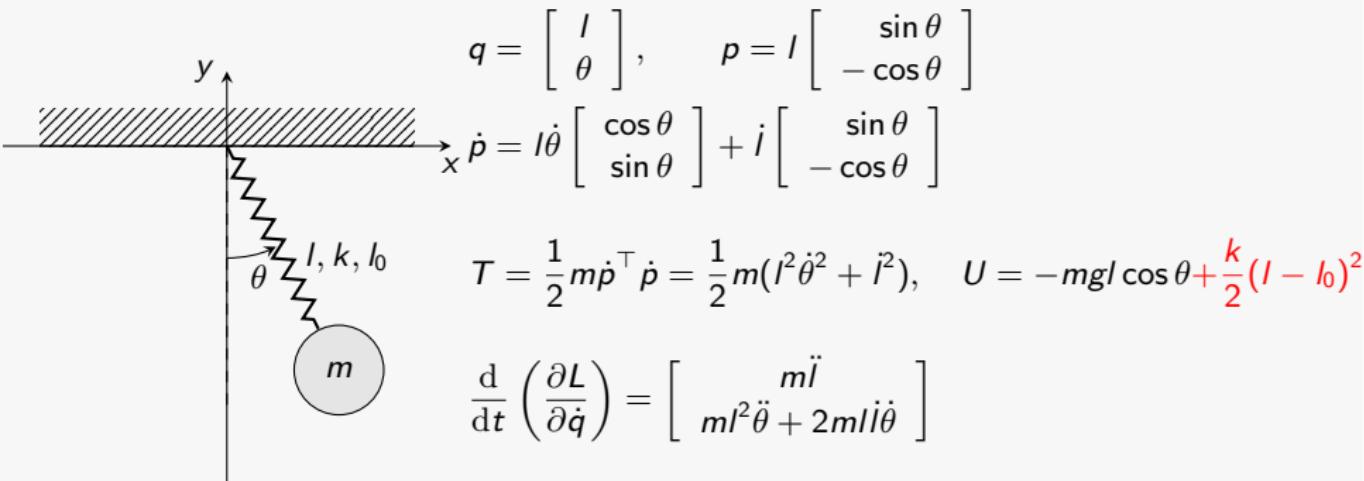
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Diagram showing an elastic pendulum of length  $l$  pivoted at the origin of a 3D coordinate system ( $x$ ,  $y$ ,  $z$  axes). The mass  $m$  is connected to the pivot by a spring with stiffness  $k$  and natural length  $l_0$ .

Position and momentum variables:

$$q = \begin{bmatrix} l \\ \theta \end{bmatrix}, \quad p = l \begin{bmatrix} \sin \theta \\ -\cos \theta \end{bmatrix}$$

$$\dot{p} = l \dot{\theta} \begin{bmatrix} \cos \theta \\ \sin \theta \end{bmatrix} + l \begin{bmatrix} \sin \theta \\ -\cos \theta \end{bmatrix}$$

Total mechanical energy:

$$T = \frac{1}{2} m \dot{p}^\top \dot{p} = \frac{1}{2} m (l^2 \dot{\theta}^2 + \dot{l}^2), \quad U = -mgl \cos \theta + \frac{k}{2} (l - l_0)^2$$

Euler-Lagrange equations:

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}} \right) = \begin{bmatrix} m \ddot{l} \\ ml^2 \ddot{\theta} + 2ml \dot{l} \dot{\theta} \end{bmatrix}$$

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Diagram of an elastic pendulum:

- A mass  $m$  is attached to a horizontal spring with stiffness  $k$  and natural length  $l_0$ .
- The spring is fixed to a wall.
- The angle  $\theta$  is measured from the vertical  $z$ -axis.
- The system is constrained to move in a plane defined by  $x$  and  $y$  axes.

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Lagrangian and potential energy:

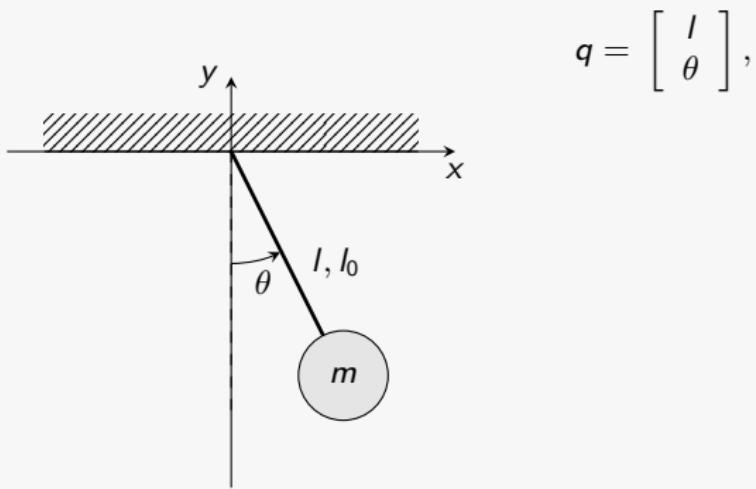
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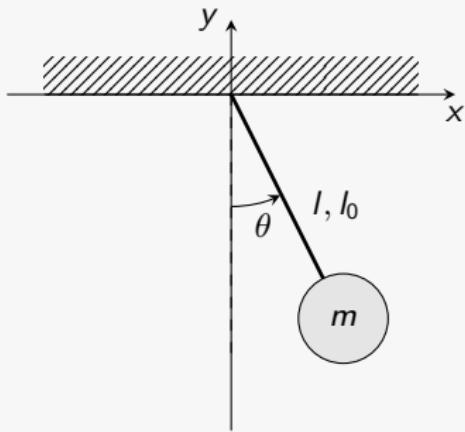
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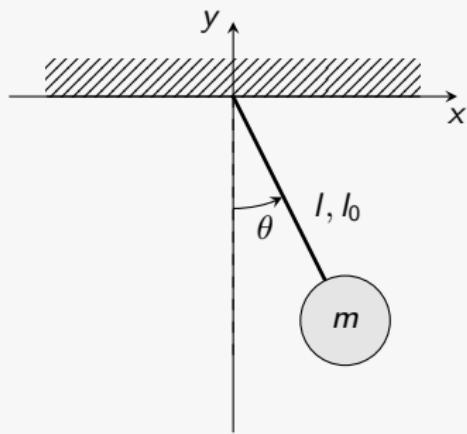
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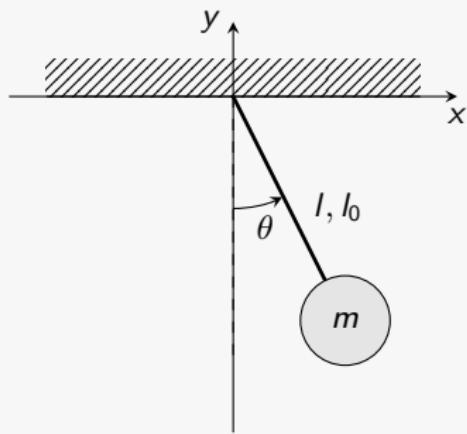
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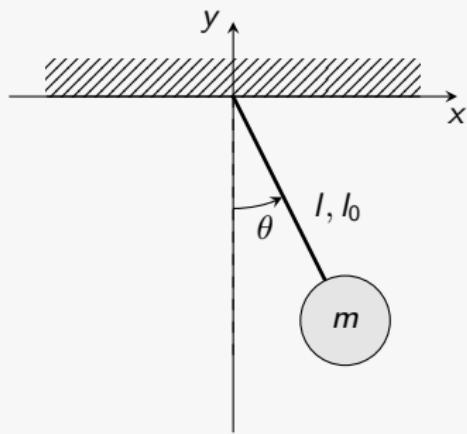
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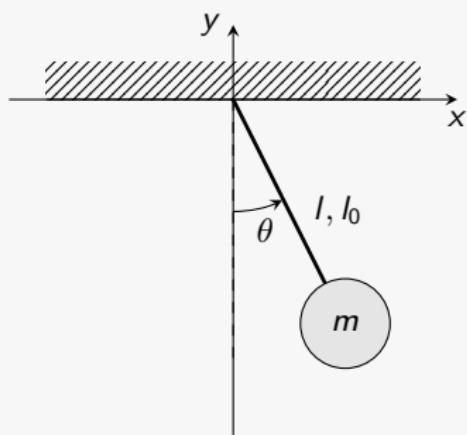
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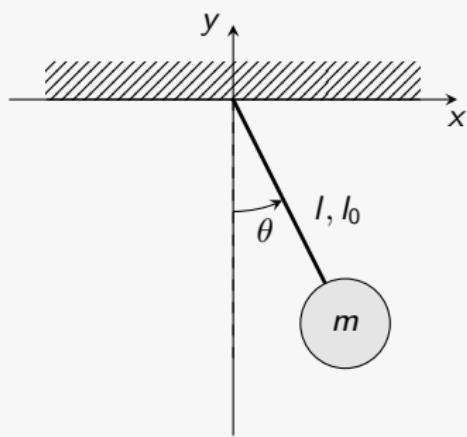
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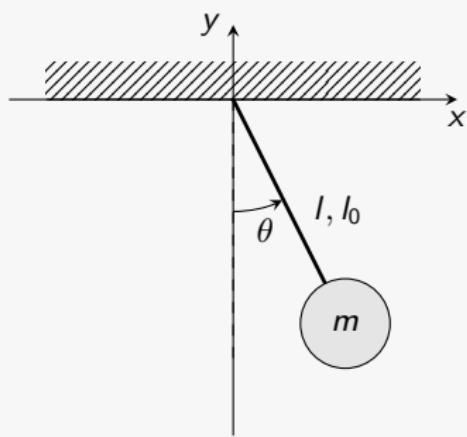
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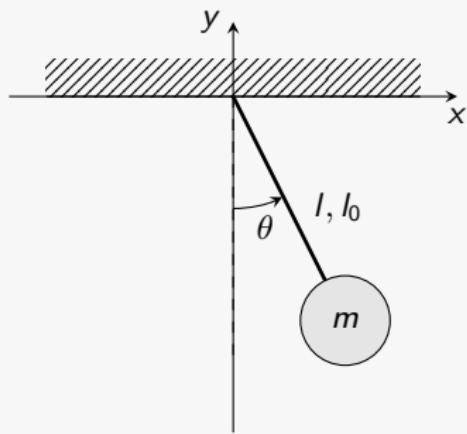
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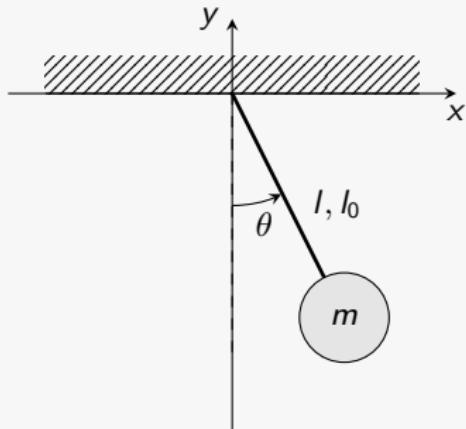
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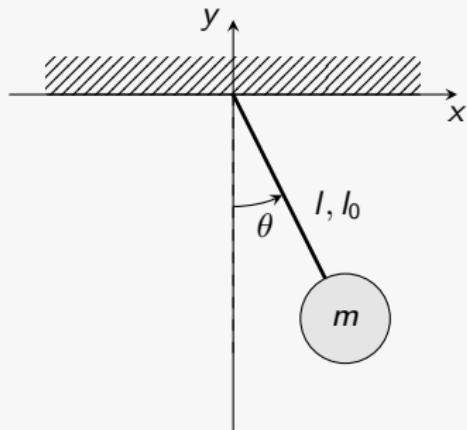
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Index-3 DAE

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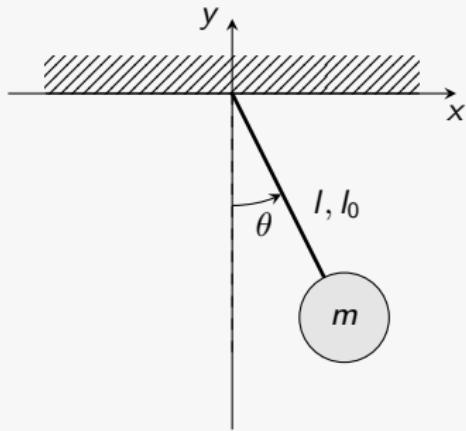


Index-3 DAE

$$\begin{bmatrix} \ddot{\theta} \\ \ddot{\theta} \\ 0 \end{bmatrix} = \begin{bmatrix} I\dot{\theta}^2 + g \cos \theta - \frac{\lambda}{m} \\ \frac{g}{l} \sin \theta - \frac{2}{l} \dot{\theta} \\ I - I_0 \end{bmatrix}$$

Index reduction:

## Example: Constrained Pendulum



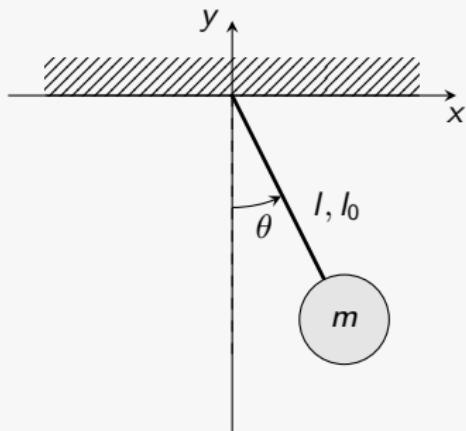
Index-3 DAE

$$\begin{bmatrix} \ddot{i} \\ \ddot{\theta} \\ 0 \end{bmatrix} = \begin{bmatrix} I\dot{\theta}^2 + g \cos \theta - \frac{\lambda}{m} \\ \frac{g}{l} \sin \theta - \frac{2}{l}\dot{i}\dot{\theta} \\ I - I_0 \end{bmatrix}$$

Index reduction:

$$\dot{C} = i,$$

## Example: Constrained Pendulum



Index-3 DAE

$$\begin{bmatrix} \ddot{i} \\ \ddot{\theta} \\ 0 \end{bmatrix} = \begin{bmatrix} I\dot{\theta}^2 + g \cos \theta - \frac{\lambda}{m} \\ \frac{g}{l} \sin \theta - \frac{2}{l} i \dot{\theta} \\ I - I_0 \end{bmatrix}$$

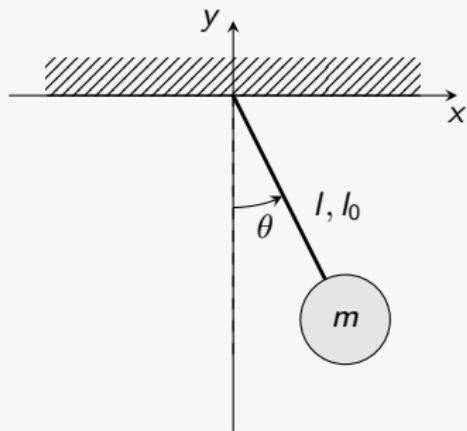
Index reduction:

$$\dot{C} = i, \quad \ddot{C} = \ddot{i}$$

Index-1 DAE

$$\begin{bmatrix} \ddot{i} \\ \ddot{\theta} \\ 0 \end{bmatrix} = \begin{bmatrix} I\dot{\theta}^2 + g \cos \theta - \frac{\lambda}{m} \\ \frac{g}{l} \sin \theta - \frac{2}{l} i \dot{\theta} \\ \ddot{i} \end{bmatrix}$$

$$C(t=0) = 0, \quad \dot{C}(t=0) = 0$$

**Example: Constrained Pendulum**

Index-3 DAE

$$\begin{bmatrix} \ddot{i} \\ \ddot{\theta} \\ 0 \end{bmatrix} = \begin{bmatrix} I\dot{\theta}^2 + g \cos \theta - \frac{\lambda}{m} \\ \frac{g}{l} \sin \theta - \frac{2}{l} i \dot{\theta} \\ I - I_0 \end{bmatrix}$$

Index reduction:

$$\dot{C} = i, \quad \ddot{C} = \ddot{i}$$

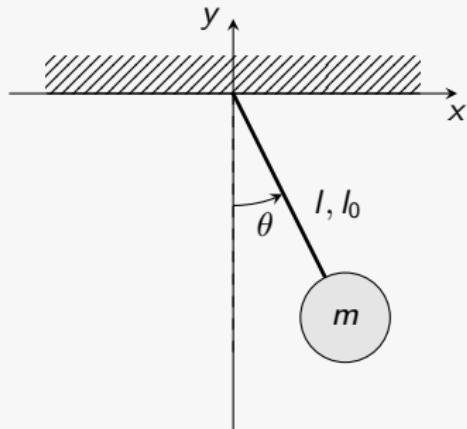
Index-1 DAE

$$\begin{bmatrix} \ddot{i} \\ \ddot{\theta} \\ 0 \end{bmatrix} = \begin{bmatrix} I\dot{\theta}^2 + g \cos \theta - \frac{\lambda}{m} \\ \frac{g}{l} \sin \theta - \frac{2}{l} i \dot{\theta} \\ \ddot{i} \end{bmatrix}$$

$$C(t=0) = 0, \quad \dot{C}(t=0) = 0$$

ODE

$$\begin{bmatrix} \ddot{\theta} \\ \lambda \end{bmatrix} = \begin{bmatrix} \frac{g}{l} \sin \theta \\ ml\dot{\theta}^2 + mg \cos \theta \end{bmatrix}, \quad I = I_0$$

**Example: Constrained Pendulum**

$\lambda$  = tension in the rod

Index-3 DAE

$$\begin{bmatrix} \ddot{i} \\ \ddot{\theta} \\ 0 \end{bmatrix} = \begin{bmatrix} I\dot{\theta}^2 + g \cos \theta - \frac{\lambda}{m} \\ \frac{g}{l} \sin \theta - \frac{2}{l} i \dot{\theta} \\ I - I_0 \end{bmatrix}$$

Index reduction:

$$\dot{C} = i, \quad \ddot{C} = \ddot{i}$$

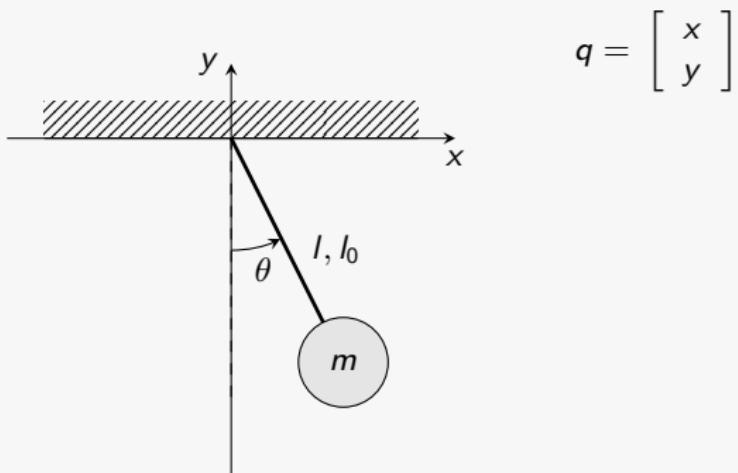
Index-1 DAE

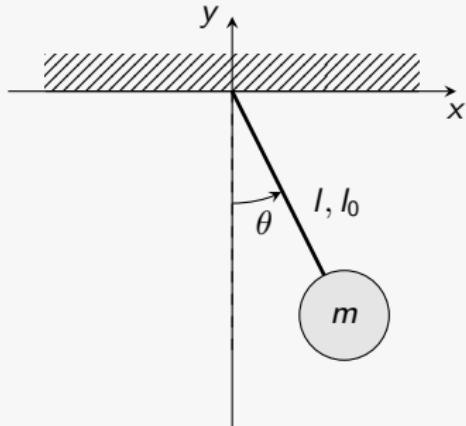
$$\begin{bmatrix} \ddot{i} \\ \ddot{\theta} \\ 0 \end{bmatrix} = \begin{bmatrix} I\dot{\theta}^2 + g \cos \theta - \frac{\lambda}{m} \\ \frac{g}{l} \sin \theta - \frac{2}{l} i \dot{\theta} \\ \ddot{i} \end{bmatrix}$$

$$C(t=0) = 0, \quad \dot{C}(t=0) = 0$$

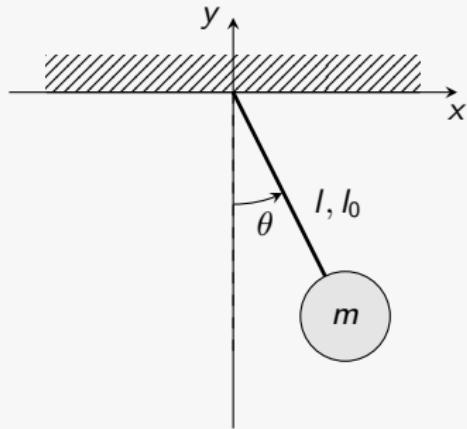
ODE

$$\begin{bmatrix} \ddot{\theta} \\ \lambda \end{bmatrix} = \begin{bmatrix} \frac{g}{l} \sin \theta \\ ml\dot{\theta}^2 + mg \cos \theta \end{bmatrix}, \quad I = I_0$$

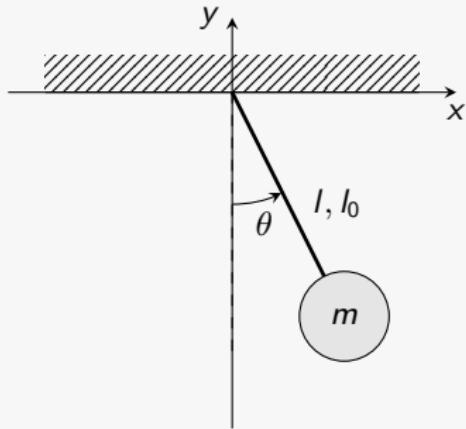
**Example: Pendulum in Cartesian Coordinates**

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$$\begin{aligned} q &= \begin{bmatrix} x \\ y \end{bmatrix} \\ p &= \begin{bmatrix} x \\ y \end{bmatrix}, \end{aligned}$$

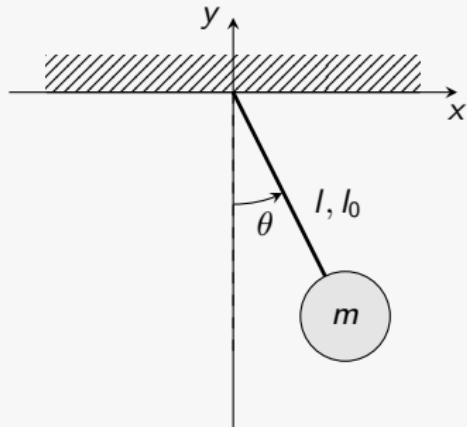
**Example: Pendulum in Cartesian Coordinates**

$$\begin{aligned} q &= \begin{bmatrix} x \\ y \end{bmatrix} \\ p &= \begin{bmatrix} x \\ y \end{bmatrix}, \quad \dot{p} = \begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} \end{aligned}$$

**Example: Pendulum in Cartesian Coordinates**

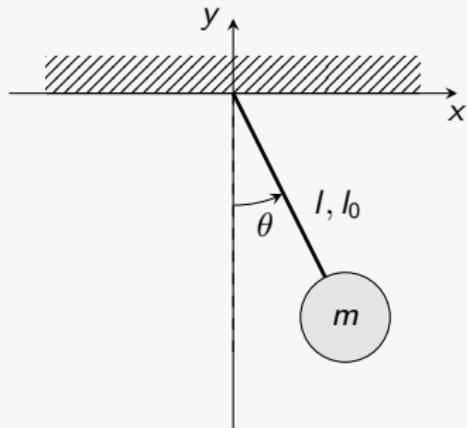
$$q = \begin{bmatrix} x \\ y \end{bmatrix}, \quad p = \begin{bmatrix} x \\ y \end{bmatrix}, \quad \dot{p} = \begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix}$$

$$T = \frac{1}{2}m\dot{p}^\top \dot{p} = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2),$$

**Example: Pendulum in Cartesian Coordinates**

$$q = \begin{bmatrix} x \\ y \end{bmatrix}, \quad p = \begin{bmatrix} x \\ y \end{bmatrix}, \quad \dot{p} = \begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix}$$

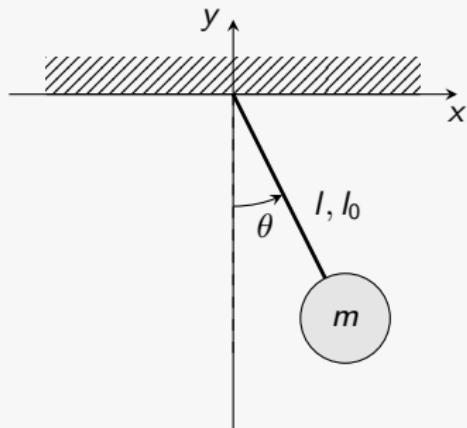
$$T = \frac{1}{2}m\dot{p}^\top \dot{p} = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2), \quad U = mgy$$

**Example: Pendulum in Cartesian Coordinates**

$$q = \begin{bmatrix} x \\ y \end{bmatrix}, \quad p = \begin{bmatrix} x \\ y \end{bmatrix}, \quad \dot{p} = \begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix}$$

$$T = \frac{1}{2}m\dot{p}^\top \dot{p} = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2), \quad U = mgy$$

$$C = \frac{1}{2}(x^2 + y^2 - l_0^2), \quad L = T - U - \lambda C$$

**Example: Pendulum in Cartesian Coordinates**

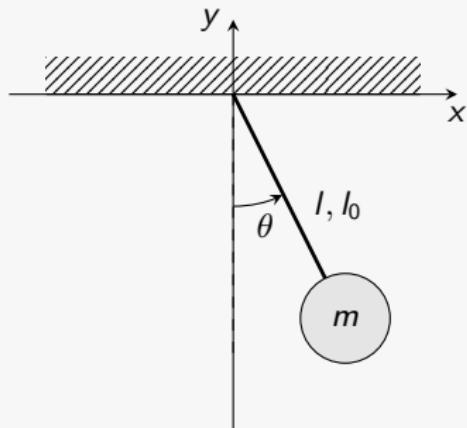
$$\begin{aligned} q &= \begin{bmatrix} x \\ y \end{bmatrix} \\ p &= \begin{bmatrix} x \\ y \end{bmatrix}, \quad \dot{p} = \begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix} \end{aligned}$$

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$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}} \right) = \begin{bmatrix} m\ddot{x} \\ m\ddot{y} \end{bmatrix}$$

## Example: Pendulum in Cartesian Coordinates



$$q = \begin{bmatrix} x \\ y \end{bmatrix}, \quad p = \begin{bmatrix} x \\ y \end{bmatrix}, \quad \dot{p} = \begin{bmatrix} \dot{x} \\ \dot{y} \end{bmatrix}$$

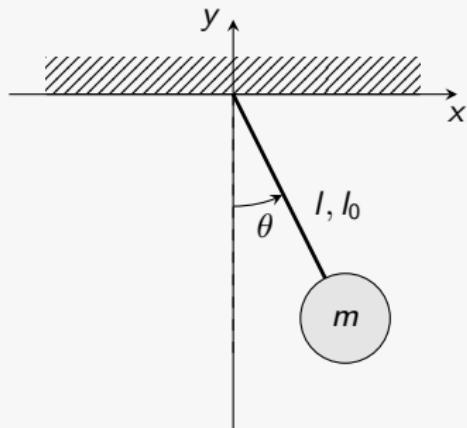
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## Example: Pendulum in Cartesian Coordinates



$$q = \begin{bmatrix} x \\ y \end{bmatrix}$$

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$$T = \frac{1}{2}m\dot{p}^\top \dot{p} = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2), \quad U = mgy$$

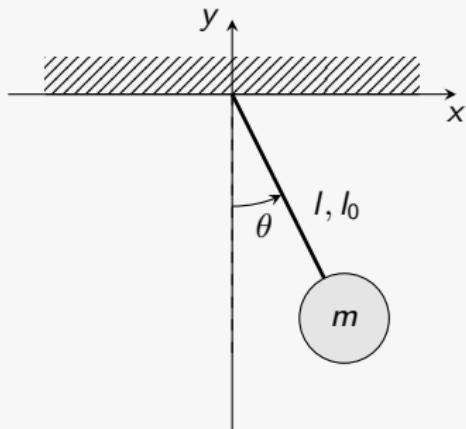
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$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ 0 \end{bmatrix} = \begin{bmatrix} -\frac{\lambda x}{m} \\ -g - \frac{\lambda y}{m} \\ \frac{1}{2}(x^2 + y^2 - l_0^2) \end{bmatrix}$$

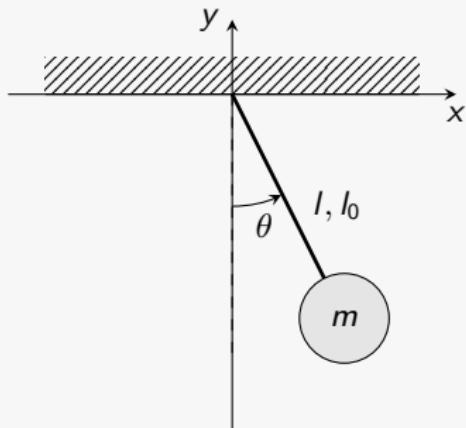
## Example: Pendulum in Cartesian Coordinates



Index-3 DAE

$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ 0 \end{bmatrix} = \begin{bmatrix} -\frac{\lambda_x}{m} \\ -g - \frac{\lambda_y}{m} \\ \frac{1}{2}(x^2 + y^2 - l_0^2) \end{bmatrix}$$

## Example: Pendulum in Cartesian Coordinates

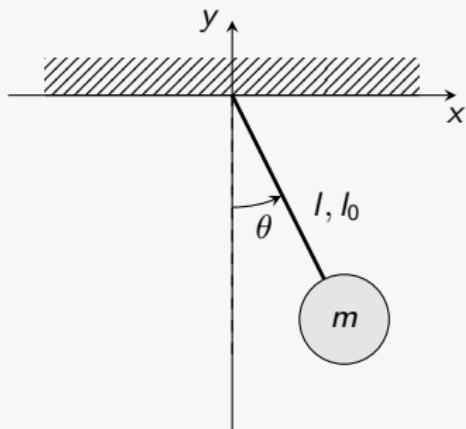


Index-3 DAE

$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ 0 \end{bmatrix} = \begin{bmatrix} -\frac{\lambda_x}{m} \\ -g - \frac{\lambda_y}{m} \\ \frac{1}{2}(x^2 + y^2 - l_0^2) \end{bmatrix}$$

Index reduction:

## Example: Pendulum in Cartesian Coordinates



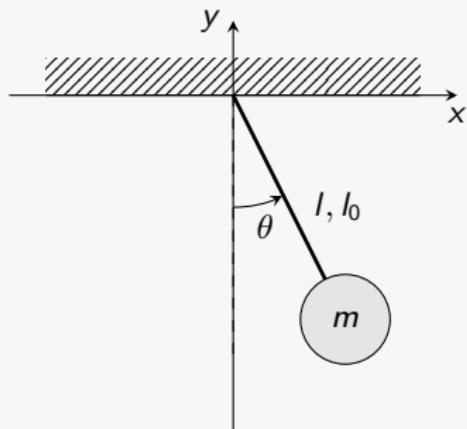
Index-3 DAE

$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ 0 \end{bmatrix} = \begin{bmatrix} -\frac{\lambda_x}{m} \\ -g - \frac{\lambda_y}{m} \\ \frac{1}{2}(x^2 + y^2 - l_0^2) \end{bmatrix}$$

Index reduction:

$$\dot{C} = \dot{x}x + \dot{y}y,$$

## Example: Pendulum in Cartesian Coordinates



Index-3 DAE

$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ 0 \end{bmatrix} = \begin{bmatrix} -\frac{\lambda_x}{m} \\ -g - \frac{\lambda_y}{m} \\ \frac{1}{2}(x^2 + y^2 - l_0^2) \end{bmatrix}$$

Index reduction:

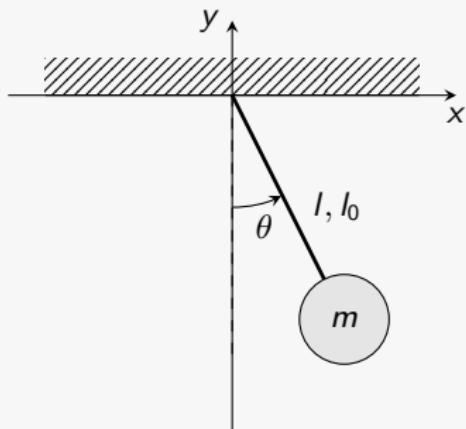
$$\dot{C} = \dot{x}x + \dot{y}y, \quad \ddot{C} = \ddot{x}x + \ddot{y}y + \dot{x}^2 + \dot{y}^2$$

Index-1 DAE

$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ 0 \end{bmatrix} = \begin{bmatrix} -\frac{\lambda_x}{m} \\ -g - \frac{\lambda_y}{m} \\ \ddot{x}x + \ddot{y}y + \dot{x}^2 + \dot{y}^2 \end{bmatrix}$$

$$C(t=0) = 0, \quad \dot{C}(t=0) = 0$$

## Example: Pendulum in Cartesian Coordinates



Index-3 DAE

$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ 0 \end{bmatrix} = \begin{bmatrix} -\frac{\lambda_x}{m} \\ -g - \frac{\lambda_y}{m} \\ \frac{1}{2}(x^2 + y^2 - l_0^2) \end{bmatrix}$$

Index reduction:

$$\dot{C} = \dot{x}x + \dot{y}y, \quad \ddot{C} = \ddot{x}x + \ddot{y}y + \dot{x}^2 + \dot{y}^2$$

Index-1 DAE

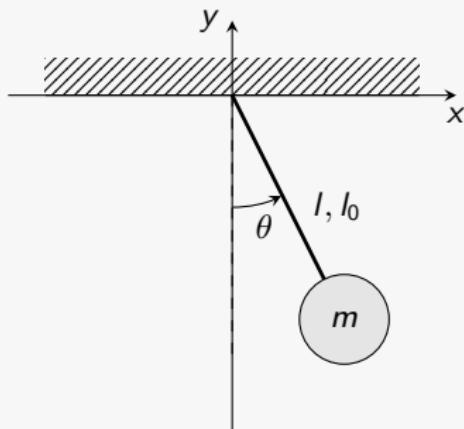
$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ 0 \end{bmatrix} = \begin{bmatrix} -\frac{\lambda_x}{m} \\ -g - \frac{\lambda_y}{m} \\ \ddot{x}x + \ddot{y}y + \dot{x}^2 + \dot{y}^2 \end{bmatrix}$$

$$C(t=0) = 0, \quad \dot{C}(t=0) = 0$$

Semi-implicit form

$$\begin{bmatrix} m & 0 & x \\ 0 & m & y \\ x & y & 0 \end{bmatrix} \begin{bmatrix} \ddot{x} \\ \ddot{y} \\ \lambda \end{bmatrix} = \begin{bmatrix} 0 \\ mg \\ \dot{x}^2 + \dot{y}^2 \end{bmatrix}$$

## Example: Pendulum in Cartesian Coordinates



$$\begin{bmatrix} \lambda x \\ \lambda y \end{bmatrix} = \text{rod tension}$$

Index-3 DAE

$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ 0 \end{bmatrix} = \begin{bmatrix} -\frac{\lambda x}{m} \\ -g - \frac{\lambda y}{m} \\ \frac{1}{2}(x^2 + y^2 - l_0^2) \end{bmatrix}$$

Index reduction:

$$\dot{C} = \dot{x}x + \dot{y}y, \quad \ddot{C} = \ddot{x}x + \ddot{y}y + \dot{x}^2 + \dot{y}^2$$

Index-1 DAE

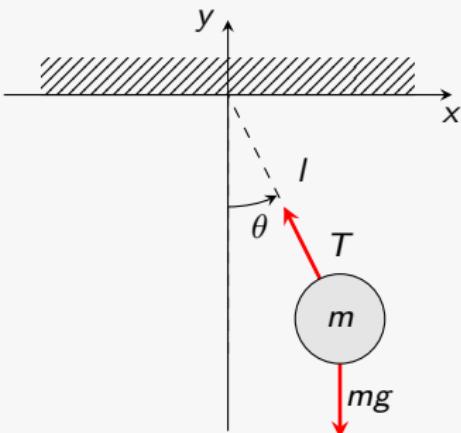
$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ 0 \end{bmatrix} = \begin{bmatrix} -\frac{\lambda x}{m} \\ -g - \frac{\lambda y}{m} \\ \ddot{x}x + \ddot{y}y + \dot{x}^2 + \dot{y}^2 \end{bmatrix}$$

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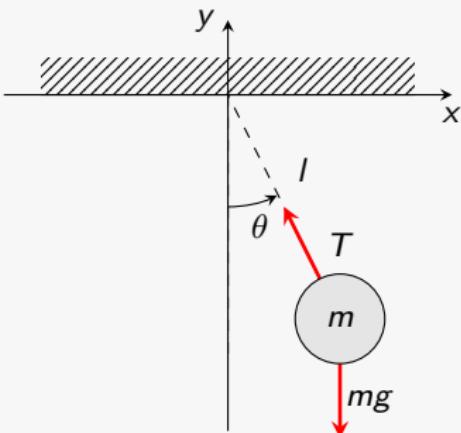
## Example: Pendulum in Cartesian Coordinates



Newton's Approach

$$m \begin{bmatrix} \ddot{x} \\ \ddot{y} \end{bmatrix} = \begin{bmatrix} T^x \\ -mg + T^y \end{bmatrix}$$

## Example: Pendulum in Cartesian Coordinates

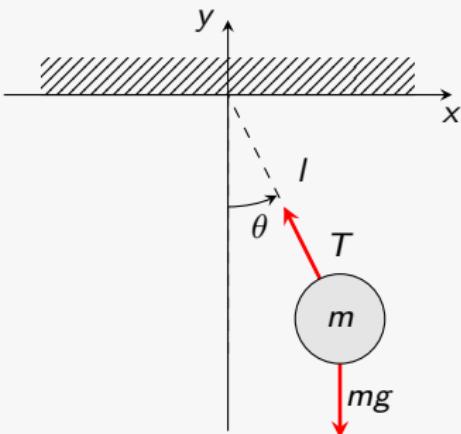


Newton's Approach

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Rod tension:

## Example: Pendulum in Cartesian Coordinates



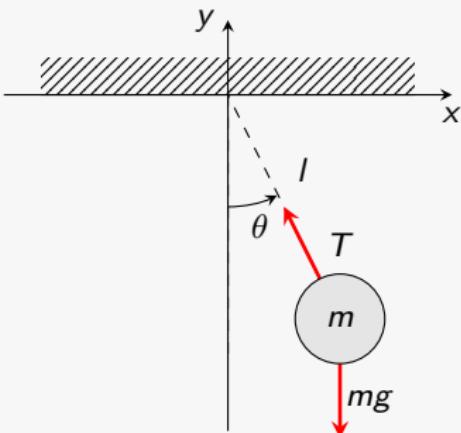
### Newton's Approach

$$m \begin{bmatrix} \ddot{x} \\ \ddot{y} \end{bmatrix} = \begin{bmatrix} T^x \\ -mg + T^y \end{bmatrix}$$

Rod tension:

$$T_g = mg \cos \theta = mg \frac{y}{\sqrt{x^2 + y^2}} = mg \frac{y}{l}$$

## Example: Pendulum in Cartesian Coordinates



### Newton's Approach

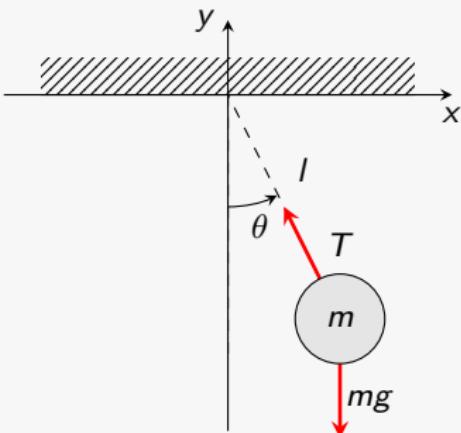
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Rod tension:

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$$T_c = ml\dot{\theta} = m \frac{\dot{x}^2 + \dot{y}^2}{l}$$

## Example: Pendulum in Cartesian Coordinates



### Newton's Approach

$$m \begin{bmatrix} \ddot{x} \\ \ddot{y} \end{bmatrix} = \begin{bmatrix} T^x \\ -mg + T^y \end{bmatrix}$$

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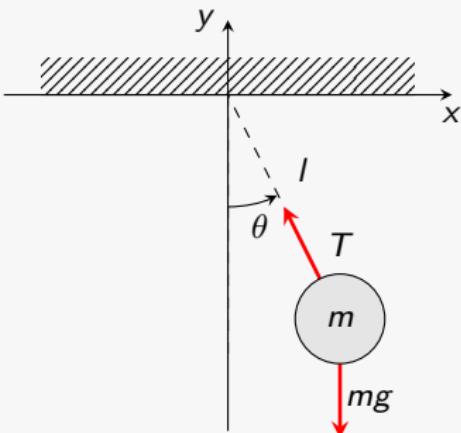
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$$T_c = ml\dot{\theta} = m \frac{\dot{x}^2 + \dot{y}^2}{l}$$

In the  $x$  and  $y$  directions:

$$T^x = T \sin \theta = T \frac{x}{\sqrt{x^2 + y^2}} = mx \frac{\dot{x}^2 + \dot{y}^2 - gy}{l^2}$$

## Example: Pendulum in Cartesian Coordinates



### Newton's Approach

$$m \begin{bmatrix} \ddot{x} \\ \ddot{y} \end{bmatrix} = \begin{bmatrix} T^x \\ -mg + T^y \end{bmatrix}$$

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$$T_g = mg \cos \theta = mg \frac{y}{\sqrt{x^2 + y^2}} = mg \frac{y}{l}$$

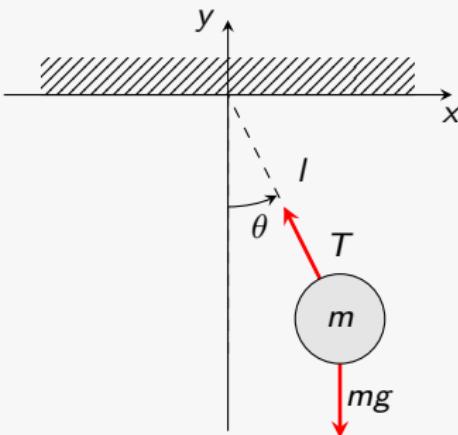
$$T_c = ml\dot{\theta} = m \frac{\dot{x}^2 + \dot{y}^2}{l}$$

In the  $x$  and  $y$  directions:

$$T^x = T \sin \theta = T \frac{x}{\sqrt{x^2 + y^2}} = mx \frac{\dot{x}^2 + \dot{y}^2 - gy}{l^2}$$

$$T^y = T \cos \theta = T \frac{y}{\sqrt{x^2 + y^2}} = my \frac{\dot{x}^2 + \dot{y}^2 - gy}{l^2}$$

## Example: Pendulum in Cartesian Coordinates



### Newton's Approach

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Rod tension:

$$T_g = mg \cos \theta = mg \frac{y}{\sqrt{x^2 + y^2}} = mg \frac{y}{l}$$

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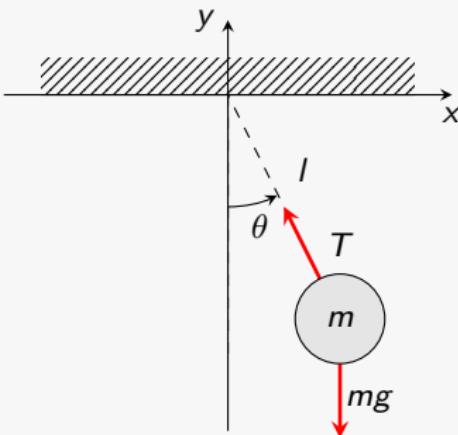
In the  $x$  and  $y$  directions:

$$T^x = T \sin \theta = T \frac{x}{\sqrt{x^2 + y^2}} = mx \frac{\dot{x}^2 + \dot{y}^2 - gy}{l^2}$$

$$T^y = T \cos \theta = T \frac{y}{\sqrt{x^2 + y^2}} = my \frac{\dot{x}^2 + \dot{y}^2 - gy}{l^2}$$

$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ 0 \end{bmatrix} = \begin{bmatrix} -\frac{\lambda_x}{m} \\ -g - \frac{\lambda_y}{m} \\ \ddot{x}x + \ddot{y}y + \frac{\dot{x}^2 + \dot{y}^2}{l^2} \end{bmatrix}$$

## Example: Pendulum in Cartesian Coordinates



### Newton's Approach

$$m \begin{bmatrix} \ddot{x} \\ \ddot{y} \end{bmatrix} = \begin{bmatrix} T^x \\ -mg + T^y \end{bmatrix}$$

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$$T_c = ml\dot{\theta} = m \frac{\dot{x}^2 + \dot{y}^2}{l}$$

In the x and y directions:

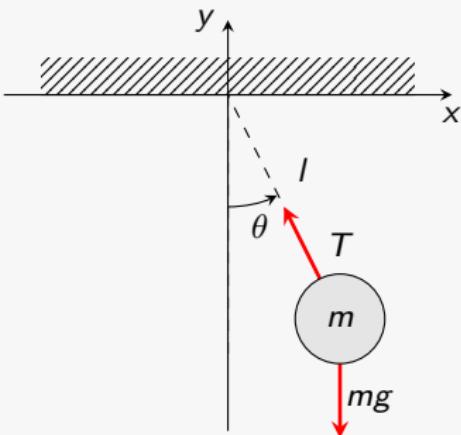
$$T^x = T \sin \theta = T \frac{x}{\sqrt{x^2 + y^2}} = mx \frac{\dot{x}^2 + \dot{y}^2 - gy}{l^2}$$

$$T^y = T \cos \theta = T \frac{y}{\sqrt{x^2 + y^2}} = my \frac{\dot{x}^2 + \dot{y}^2 - gy}{l^2}$$

$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ 0 \end{bmatrix} = \begin{bmatrix} -\frac{\lambda x}{m} \\ -g - \frac{\lambda y}{m} \\ \ddot{x}x + \ddot{y}y + \dot{x}^2 + \dot{y}^2 \end{bmatrix}$$

$$-\lambda x^2 - mgy - \lambda y^2 + m\dot{x}^2 + m\dot{y}^2$$

## Example: Pendulum in Cartesian Coordinates



### Newton's Approach

$$m \begin{bmatrix} \ddot{x} \\ \ddot{y} \end{bmatrix} = \begin{bmatrix} T^x \\ -mg + T^y \end{bmatrix}$$

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$$T^x = T \sin \theta = T \frac{x}{\sqrt{x^2 + y^2}} = mx \frac{\dot{x}^2 + \dot{y}^2 - gy}{l^2}$$

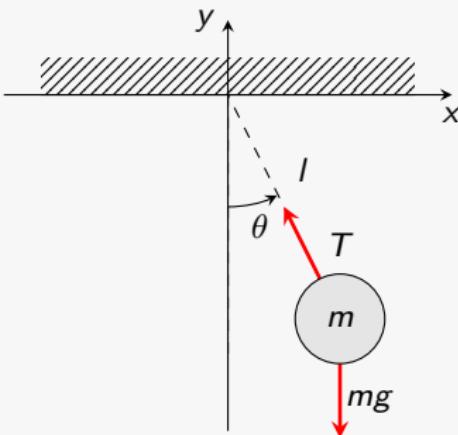
$$T^y = T \cos \theta = T \frac{y}{\sqrt{x^2 + y^2}} = my \frac{\dot{x}^2 + \dot{y}^2 - gy}{l^2}$$

$$\begin{bmatrix} \ddot{x} \\ \ddot{y} \\ 0 \end{bmatrix} = \begin{bmatrix} -\frac{\lambda x}{m} \\ -g - \frac{\lambda y}{m} \\ \ddot{x}x + \ddot{y}y + \dot{x}^2 + \dot{y}^2 \end{bmatrix}$$

$$-\lambda x^2 - mgy - \lambda y^2 + m\dot{x}^2 + m\dot{y}^2$$

$$\lambda = m \frac{\dot{x}^2 + \dot{y}^2 - gy}{x^2 + y^2} = m \frac{\dot{x}^2 + \dot{y}^2 - gy}{l^2}$$

## Example: Pendulum in Cartesian Coordinates



### Newton's Approach

$$m \begin{bmatrix} \ddot{x} \\ \ddot{y} \end{bmatrix} = \begin{bmatrix} T^x \\ -mg + T^y \end{bmatrix}$$

Rod tension:

$$T_g = mg \cos \theta = mg \frac{y}{\sqrt{x^2 + y^2}} = mg \frac{y}{l}$$

$$T_c = ml\dot{\theta} = m \frac{\dot{x}^2 + \dot{y}^2}{l}$$

In the  $x$  and  $y$  directions:

$$T^x = T \sin \theta = T \frac{x}{\sqrt{x^2 + y^2}} = mx \frac{\dot{x}^2 + \dot{y}^2 - gy}{l^2}$$

$$T^y = T \cos \theta = T \frac{y}{\sqrt{x^2 + y^2}} = my \frac{\dot{x}^2 + \dot{y}^2 - gy}{l^2}$$

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$$\lambda = m \frac{\dot{x}^2 + \dot{y}^2 - gy}{x^2 + y^2} = m \frac{\dot{x}^2 + \dot{y}^2 - gy}{l^2}$$

$$T^x = \lambda x, \quad T^y = \lambda y$$

## Principle of Virtual Works

How to include external forces/torques?

- Generalised coordinates  $q \Rightarrow$  generalised forces  $F^q$

Typically, it is natural to express forces  $F$  in a coordinate frame  $x \neq q$ .

How can we compute  $F^q$ ?

- Virtual displacement:  $\delta x$
- Principle of Virtual Works:  $\delta W := F\delta x = F^q\delta q$
- $x = x(q)$ , then:  $\delta x = \frac{dx}{dq}\delta q$

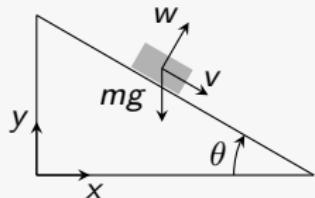
Generalised forces:

$$F^q = F \frac{dx}{dq}$$

## Examples

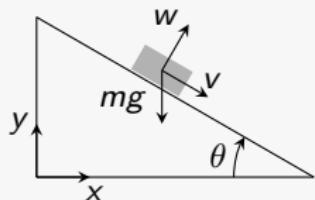
## Examples

Mass on a frictionless inclined plane



## Examples

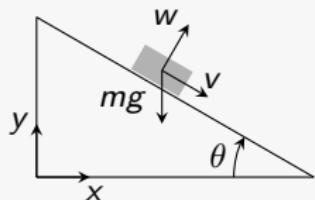
Mass on a frictionless inclined plane



$$\begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} v \cos \theta + w \sin \theta \\ -v \sin \theta + w \cos \theta \end{bmatrix} = \underbrace{\begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix}}_R \begin{bmatrix} v \\ w \end{bmatrix}$$

## Examples

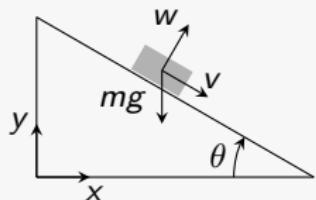
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$$q = v,$$

## Examples

Mass on a frictionless inclined plane

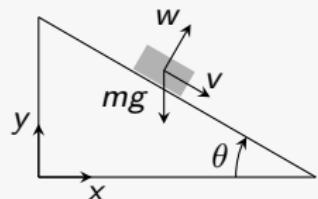


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$$q = v, \quad F\delta y = F^q\delta v,$$

## Examples

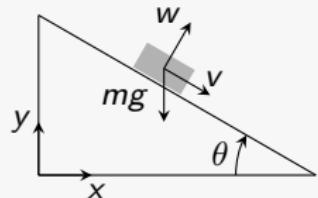
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$$q = v, \quad F\delta y = F^q \delta v, \quad F^q = F \frac{dy}{dv} = -F \sin \theta = mg \sin \theta$$

## Examples

Mass on a frictionless inclined plane



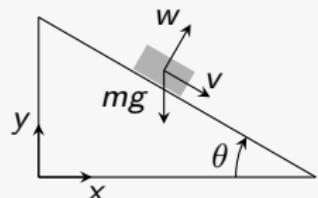
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$$q = v, \quad F\delta y = F^q \delta v, \quad F^q = F \frac{dy}{dv} = -F \sin \theta = mg \sin \theta$$

The generalised force  $F^q$  is constant!

## Examples

Mass on a frictionless inclined plane

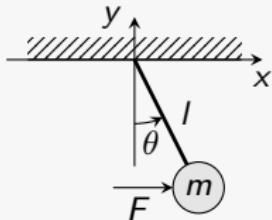


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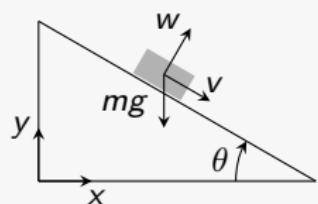
**The generalised force  $F^q$  is constant!**

Pendulum with Horizontal Force



## Examples

Mass on a frictionless inclined plane

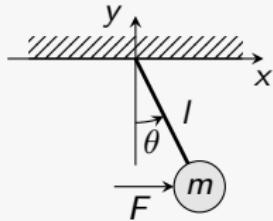


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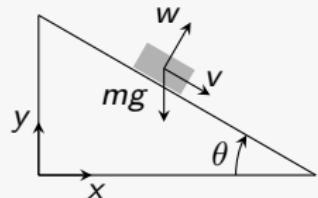
Pendulum with Horizontal Force



$$q = \theta,$$

## Examples

Mass on a frictionless inclined plane

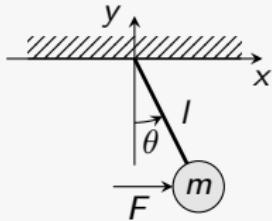


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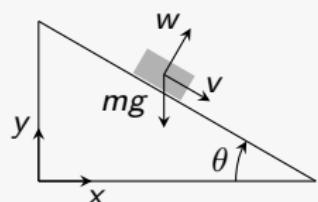
Pendulum with Horizontal Force



$$q = \theta, \quad x = l \sin \theta,$$

## Examples

Mass on a frictionless inclined plane

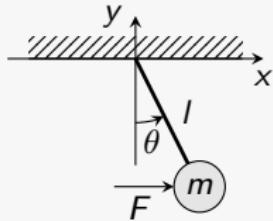


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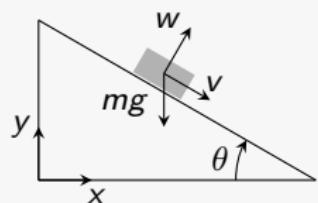
Pendulum with Horizontal Force



$$q = \theta, \quad x = l \sin \theta, \quad F^q = F \frac{dx}{d\theta} = Fl \cos \theta$$

## Examples

Mass on a frictionless inclined plane

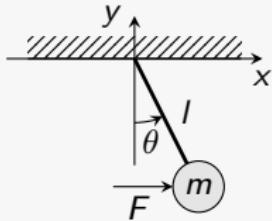


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Pendulum with Horizontal Force

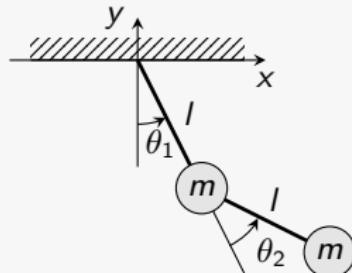
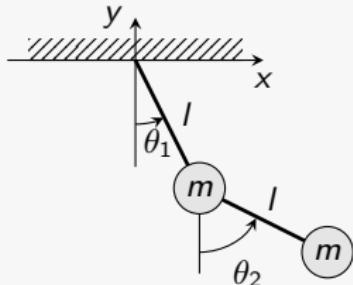


$$q = \theta, \quad x = l \sin \theta, \quad F^q = F \frac{dx}{d\theta} = Fl \cos \theta$$

The generalised force  $F^q$  depends on  $q$ !

## Exercises

- Model a double pendulum in polar and cartesian coordinates (assume  $l_1 = l_2 = l$ ,  $m_1 = m_2 = m$ )



- Model a pendulum attached to a crane in polar and cartesian coordinates
- Assume to be able to change the length of the crane rod

